## Regular article

# Core-valence correlation effects for molecules containing first-row atoms. Accurate results using effective core polarization potentials $*$

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Abstract. The accuracy of employing effective core polarization potentials (CPPs) to account for the effects of core-valence correlation on the spectroscopic constants and dissociation energies of the molecules  $B_2, C_2$ ,  $N_2$ ,  $O_2$ ,  $F_2$ ,  $CO$ ,  $CN$ ,  $CH$ ,  $HF$ , and  $C_2H_2$  has been investigated by comparison to accurate all-electron benchmark calculations. The results obtained from the calculations employing CPPs were surprisingly accurate in every case studied, reducing the errors in the calculated valence  $D<sub>e</sub>$  values from a maximum of nearly  $2.5$  kcal/mol to just 0.3 kcal/mol. The effects of enlarging the basis set and using higher-order valence electron correlation treatments were found to have only a small influence on the core-valence correlation effect predicted by the CPPs. Thus, to accurately recover the effects of intershell correlation, effective core polarization potentials such as the ones used in the present work provide an attractive alternative to carrying out computationally demanding calculations where the core electrons are explicitly included in the correlation treatment.

Key words: Core-valence correlation  $-$ Core polarization potentials

#### 1 Introduction

One of the most common assumptions in quantum chemistry involves the intuitive notion that molecular structure and energetics are primarily determined by the interactions of the valence electrons. Within this description the core electrons provide a source for the potential that is experienced by the valence electrons, but, because the core and valence electrons are concentrated in different regions of space, other effects,

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including intershell correlation effects, are negligible. This means that a simple orbital description would suffice for the description of the core electrons.

In 1975, Meyer and Rosmus [1] reported the first study that systematically investigated the influence of core-valence correlation on molecular properties in ab initio calculations. Their work on the first- and second-row diatomic hydrides, as well as the LiNa dimer [2], demonstrated that core-valence correlation can be nearly as important as valence correlation for alkali and alkaline earth compounds. Even for compounds containing the atoms B to F, if the goal of the calculation is the prediction of very accurate molecular structures (within thousandths of Angstroms) or energetics (better than 1 kcal/mol), then the effects of correlating the lowlying core electrons must be taken into account (see, for example, [3-6]). This dramatically increases the computational effort because more electron pairs must be correlated and larger basis sets must be used to simultaneously describe valence, core-valence, and core correlation effects. Of course, this level of effort is only worthwhile if the errors introduced by finite one-particle basis sets and the valence correlation treatment are smaller than the core-valence correlation effects.

A careful analysis of core-valence correlation reveals that the physical origin of this effect is principally due to the dynamic polarization of the atomic cores by the valence electrons [7]. This suggests that such effects could be taken into account by means of an effective core polarization potential (CPP). In fact this has been done for decades to improve model potentials [8–11]. However, it was again Meyer and his group who in the early 1980 s incorporated an effective core polarization potential into ab initio calculations  $[12–14]$ . They emphasized the importance of not only including the oneelectron terms in such a potential, but also the proper implemention of the two-electron terms. Furthermore, they discussed different functional forms to cut-off the CPP at short distances, thus accounting for valence electrons that penetrate the cores in the determination of the polarizing field. Meyer and co-workers [12] found that the performance of the different cut-off functions

<sup>\*</sup>Dedicated to Prof. Dr. Wilfried Meyer on the occasion of his  $60<sup>th</sup>$ birthday

was almost identical as long as the one adjustable parameter was fitted to experimental ionization potentials. This adjustment procedure was claimed to be a simple means to account for higher-order polarizabilities, as well as nonadiabatic and exclusion effects. With these CPPs they reported highly accurate spectroscopic constants for alkali and alkaline earth diatomics [13]. On the other hand, they also determined that the core-valence effect was overestimated for  $d$ -electrons. The latter problem was addressed by Leininger et al. [15] by introducing a nonlocal cut-off function.

Since the work of Meyer and co-workers, only a few investigations have used effective core polarization potentials within an all-electron model  $[16–19]$ . On the other hand, effective core polarization potentials have become popular in connection with pseudopotentials [20–38]. Within the pseudopotential scheme, CPPs not only account for core-valence correlation effects, but also describe the static polarization of the atomic cores by the molecular environment, an effect that is missing in standard pseudopotentials. Despite their wide use, only a few attempts have been made to compare results from such calculations with respective calculations in which the core electrons were explicitly correlated [33, 39, 40]. The findings in these investigations were always that the CPP results and those from explicitly correlating the core electrons were in reasonable correspondence, and the authors generally claimed that the agreement would probably be even better if more accurate all-electron calculations were possible.

Thus, there is considerable interest in comparing high accuracy calculations of core-valence correlation effects with similar calculations applying an effective core polarization potential. In previous benchmark studies  $[6, 41-43]$ , systematic sequences of correlation consistent core-valence basis sets were used to predict the complete basis set limits for the effects of 1s correlation on the spectroscopic constants of several diatomics containing first-row atoms. In the present study these very accurate results, as well as other results from the literature [5, 44– 47], have been used to benchmark the utility of effective CPPs in recovering the effects of intershell correlation in molecules containing first-row atoms. The motivation for this work is strongly driven by the high computational expense associated with carrying out accurate calculations correlating all of the electrons, especially in the case of polyatomic molecules. As will be demonstrated in the present study, the use of effective core polarization potentials appear to be capable of accurately reproducing core-valence correlation effects with minimal computational cost.

#### 2 Computational methodology

All the results reported in this paper were obtained in all-electron calculations using the coupled-cluster method with single and double substitutions (CCSD) [48, 49] and CCSD with a perturbative inclusion of triples [CCSD(T)] [50, 51]. The explicit correlation treatment was restricted to only the valence electrons, i.e., the frozen core approximation was used throughout. To describe corevalence correlation effects, we supplemented the molecular Hamiltonian with an effective core-polarization potential [12]:

$$
V_{\text{CPP}} = -\frac{1}{2} \sum_{\lambda} \alpha_{\lambda} \mathbf{f}_{\lambda}^{2}
$$

$$
\mathbf{f}_{\lambda} = \sum_{\lambda} \frac{\mathbf{r}_{\lambda i}}{r_{\lambda i}^{3}} g_{\lambda}(r_{\lambda i}) - \sum_{\mu(\neq \lambda)} \frac{Q_{\mu} \mathbf{r}_{\lambda \mu}}{r_{\lambda \mu}^{3}} g_{\lambda}(r_{\lambda \mu})
$$

$$
g_{\lambda}(r) = [1 - \exp(-\gamma_{\lambda n} r^{2})]^{n}
$$

The summation is over all polarizable cores  $\lambda$  in the molecule with the dipole polarizabilities  $\alpha_{\lambda}$  and the electric field  $f_{\lambda}$  at the site of core  $\lambda$  generated by the valence electrons i and the other cores  $\mu$ . The function  $g_{\lambda}(r)$  cuts off the polarizing field inside the core region and thus prevents the CPP from becoming singular. We tested two forms of cut-off functions differing in the exponential parameter  $n$ . Calculations using  $n = 1$  are subsequently denoted CPP/1, while calculations with  $n = 2$  are labeled CPP/2. Since the core-polarization potential defined in this way accounts not only for corevalence correlation effects but also for the static polarization of the atomic cores in the molecular environment, we forced our cores to be spherical by taking the 1s orbitals from calculations on the atomic ground states and freezing these in the molecular calculations with the CPP. For the reference calculations without the CPP, the cores were allowed to relax. Thus, the difference between the two calculations is exactly the core-valence correlation contribution. The dipole polarizabilities for the 1s cores were taken from coupled Dirac-Fock calculations [52]. The parameters  $\gamma_{\lambda n}$  were adjusted as suggested by Meyer and co-workers [12] such that corevalence correlation contributions to the ionization potentials of single valence electron ions derived from experimental ionization potentials could be reproduced [53]. The CPP parameters are displayed in Table 1.

The basis sets used in the present work were those of Dunning and co-workers [54, 55], namely the correlation consistent polarized valence double-zeta through quadruple-zeta sets (cc-pVnZ,  $n = D$ ,  $T$ , Q) and the respective sets augmented by a set of diffuse functions (aug-cc-pV $nZ$ ). For CO we also tested the impact of special corevalence correlation functions (cc-pCVnZ) [56].

Spectroscopic constants for the diatomic molecules considered in this study  $(B_2, C_2, N_2, O_2, F_2, CO, CN, CH, and HF)$  were determined by first fitting 7-9 energy points to polynomials in  $\Delta r = r-r_e$  and then using the resulting polynomial coefficients to obtain minimum energies, equilibrium bond lengths, and harmonic frequencies. Results for other spectroscopic constants, e.g., vibration-rotation interaction and anharmonicity constants, are available on request from the authors. The equilibrium geometry and minimum energies of  $C_2H_2$  were obtained by fitting a total of 13 energies to polynomials in  $\Delta r$ (CH) and  $\Delta r$ (CC). In the atomic calculations required for the determination of the dissociation energies, full symmetry equivalencing was used on the molecular orbitals. Lastly, in the case of CH, the SCF orbitals were state averaged to produce the correct  ${}^{2}\Pi$  symmetry (of course, only one component was used in the subsequent correlated calculations). For the atomic and open-shell molecules, the spin-restricted RCCSD(T) method of Knowles et al. [57] was employed. All calculations were performed with the program package MOLPRO.<sup>1</sup>

#### 3 Results and discussion

The 1s core contributions to  $D_e$ ,  $r_e$ , and  $\omega_e$  ( $\Delta D_e$ ,  $\Delta r_e$ , and  $\Delta\omega_e$ ) are shown in Tables 2–11 for B<sub>2</sub>, C<sub>2</sub>, N<sub>2</sub>, O<sub>2</sub>, F<sub>2</sub>, CO, CN, CH, HF, and  $C_2H_2$ , respectively. In each case the core-valence correlation effects on  $D_e$ ,  $r_e$ , and  $\omega_e$ 

<sup>&</sup>lt;sup>1</sup> MOLPRO is a package of ab initio programs written by H.-J Werner and P. J Knowles with contributions from J. Almlöf, R. D. Amos, A. Berning, D. L. Cooper, M. J. O. Deegan, A. J. Dobbyn, F. Eckert, S. T. Elbert, C. Hampel, R. Lindh, A. W. Lloyd, W. Meyer, M. E. Mura, A. Nicklass, K. A. Peterson, R. M. Pitzer, P. Pulay, M. Schütz, H. Stoll, A. J. Stone, P. R. Taylor, and T. Thorsteinsson

**Table 1.** Parameters  $\alpha$  and  $\gamma$  for the core polarization potentials of B through F

	α	$\gamma$ ( $n = 1$ )	$\gamma$ (n = 2)
$B^{3+}$	0.01953	1.8942	4.2385
$C^{4+}$	0.008919	2.8011	6.2557
$\overline{N}^{5+}$	0.004633	3.8296	8.4652
$O^{6+}$	0.002639	4.9942	10.9309
$\mathbf{F}^{7+}$	0.001611	6.1769	13.2663

obtained with the CPPs are given for SCF, CCSD, and  $CCSD(T)$  wave functions with the cc-pVDZ, cc-pVTZ, and cc-pVQZ basis sets. The CPP results are compared in these tables to the best estimates for these quantities obtained from high quality calculations reported in the literature, where all the electrons were included in the correlation treatment and larg e core-valence basis sets were used. In most cases these involve d CCSD(T) calculations calibrate d by internally contracted multireference averaged coupled pair functional (ACPF) [58] results. In three cases, CN, CH, an d HF, new ACPF calculations with the cc-pCV5Z basis set were carried out to provide more accurate estimates of the corevalence correlation effects. The results obtained in these calculations have been adopted as our "best estimates" for these species. Figure 1 summarizes our present results by comparing the CPP/1 and CPP/2 1s corevalence contributions,  $\Delta D_e$ ,  $\Delta r_e$ , and  $\Delta \omega_e$ , obtained with the cc-pVQZ basis sets at the CCSD(T) level of theory to the best estimate values for these quantities.

## 3.1 Homonuclear diatomics:  $B_2$ ,  $C_2$ ,  $N_2$ ,  $O_2$ ,  $F_2$

The contributions to  $D_e$ ,  $r_e$ , and  $\omega_e$  from correlating the 1s-like core electrons in the first-row homonuclear diatomics are shown in Tables 2–6. These results demonstrate two very genera l trend s in regards to the core-valence correlation effects predicted by the present CPPs. Firstly, the effect of valence electron correlation, while naturally having a very larg e impact on the total value s of the spectroscopic constants, has only a relatively small effect on the 1s core-valence contribution to these properties as predicted by either CPP/ 1 or CPP/ 2. For example, in the case of the dissociation energy of  $B_2$ , the CPP/2 result for  $\Delta D_e$  using an SCF wave function and the cc-pVQZ basis set is 0.42 kcal/mol, whil e using a CCSD(T) correlation treatment for the valence electrons yields a value of  $\Delta D_e$  of 0.96 kcal/mol. This is to be compared to an  $SCF$ /cc-pVQZ  $D_e$  of just 21.1 kcal/mol and a CCSD(T)/cc-pVQZ valu e of 64.1 kcal/mol. In addition, the results for  $\Delta D_{\rm e}$  of  ${\rm F_2}$  are also very similar at the SCF and CCSD(T) levels  $(-0.04 \text{ and } -0.01 \text{ kcal})$ mol), despite the fact that the molecule is not even

Fig. 1a–c. The 1s core-valence contributions to a  $D_e$ , b  $r_e$ , and c  $\omega_e$ obtained using effective core polarization potentials at the CCSD(T)/cc-pVQZ level of theory are compared to the best estimates obtained in large basis set, all-electrons correlated calculations. The CPP $/1$  results are designated by the *open symbols*, while the CPP/2 results are shown with  $\hat{m}$ lled symbols

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Method	Basis set	$D_e$ (kcal/mol)			$r_e$ (Å)			$\omega_e$ (cm <sup>-1</sup> )		
		Valence	CPP/1	CPP/2	Valence	CPP/1	CPP/2	Valence	CPP/1	CPP/2
$CCSD(T)/cc$ -p $CV6Z^a$		64.77		$+0.76$	1.5919		$-0.0064$	1050.4		$+8.9$
Best estimate <sup>a</sup>			$+1.1$			$-0.007$			$+10$	
<b>SCF</b>	$cc$ -p $VDZ$ $cc$ -p $VTZ$ $cc$ -p $VQZ$	19.52 20.84 21.07	0.27 0.29 0.30	0.38 0.40 0.42	1.6500 1.6388 1.6364	$-0.0062$ $-0.0066$ $-0.0066$	$-0.0077$ $-0.0083$ $-0.0084$	940.7 936.7 937.0	7.1 7.3 7.4	10.1 10.3 10.6
<b>CCSD</b>	$cc$ -p $VDZ$ $cc$ -p $VTZ$ $cc-pVQZ$	49.95 53.43 54.65	0.56 0.63 0.64	0.83 0.92 0.95	1.6290 1.6031 1.5972	$-0.0056$ $-0.0058$ $-0.0058$	$-0.0070$ $-0.0073$ $-0.0074$	1009.3 1030.6 1037.0	6.9 7.3 7.4	9.9 10.4 10.6
CCSD(T)	$cc$ -p $VDZ$ $cc$ -p $VTZ$ $cc-pVQZ$	57.74 62.51 64.14	0.53 0.60 0.61	0.83 0.93 0.96	1.6272 1.5999 1.5936	$-0.0055$ $-0.0057$ $-0.0056$	$-0.0069$ $-0.0072$ $-0.0073$	1020.3 1042.1 1048.6	6.6 7.0 7.0	9.5 10.0 10.1

Table 2. Calculated spectroscopic constants of  $B_2$  for valence-only correlation and the changes in these values due to using core polarization potentials

 $^a$  Ref. [6]

Table 3. Calculated spectroscopic constants of  $C_2$  for valence-only correlation and the changes in these values due to using core polarization potentials

Method	Basis set	$D_e$ (kcal/mol)			$r_e$ (Å)			$\omega_e$ (cm <sup>-1</sup> )		
		Valence	CPP/1	CPP/2	Valence	CPP/1	CPP/2	Valence	CPP/1	CPP/2
$CCSD(T)/cc$ -p $CV6Z^a$		144.97	$+1.00$		1.2443		$-0.0032$	1860.1		$+10.6$
Best estimate <sup>a</sup>			$+1.5$			$-0.004$			$+13$	
<b>SCF</b>	$cc$ -p $VDZ$ $cc$ -p $VTZ$ $cc$ -p $VQZ$	13.97 17.65 18.40	0.51 0.63 0.65	0.67 0.82 0.85	1.2523 1.2409 1.2389	$-0.0030$ $-0.0033$ $-0.0033$	$-0.0037$ $-0.0041$ $-0.0042$	1914.3 1904.6 1905.1	9.2 10.0 10.0	12.5 13.5 13.6
<b>CCSD</b>	$cc$ -p $VDZ$ cc-pVTZ $cc$ -p $VQZ$	113.80 121.02 124.35	0.93 1.02 1.03	1.36 1.47 1.49	1.2662 1.2469 1.2421	$-0.0028$ $-0.0030$ $-0.0030$	$-0.0034$ $-0.0038$ $-0.0038$	1861.6 1881.3 1892.6	7.8 8.5 8.6	10.8 11.7 11.9
CCSD(T)	$cc$ -p $VDZ$ $cc$ -p $VTZ$ $cc$ -p $VQZ$	130.05 139.39 143.25	0.88 0.99 0.99	1.37 1.50 1.51	1.2705 1.2507 1.2458	$-0.0028$ $-0.0030$ $-0.0030$	$-0.0035$ $-0.0038$ $-0.0038$	1828.2 1845.7 1856.3	7.6 8.1 8.3	10.6 11.2 11.5

 $^a$  Ref. [6].

Table 4. Calculated spectroscopic constants of  $N_2$  for valence-only correlation and the changes in these values due to using core polarization potentials

Method	Basis set	$D_e$ (kcal/mol)			$r_e$ (Å)			$\omega_e$ (cm <sup>-1</sup> )		
		Valence	CPP/1	CPP/2	Valence	CPP/1	CPP/2	Valence	CPP/1	CPP/2
$CCSD(T)/cc$ -p $CV6Z^a$ Best estimate <sup>a</sup>		226.49	$+0.80$ $+0.9$		1.0991	$-0.0021$ $-0.002$		2361.2	$+9.9$ $+9$	
<b>SCF</b>	$cc$ -p $VDZ$	112.15	0.70	0.92	1.0773	$-0.0018$	$-0.0021$	2758.1	7.9	10.5
	$cc$ - $pVTZ$	120.38	0.84	1.10	1.0671	$-0.0020$	$-0.0025$	2731.5	8.9	11.5
	$cc$ -p $VQZ$	121.81	0.85	1.12	1.0656	$-0.0020$	$-0.0025$	2729.5	9.0	11.8
<b>CCSD</b>	$cc-pVDZ$	193.67	0.85	1.14	1.1128	$-0.0018$	$-0.0022$	2408.9	7.2	9.8
	$cc$ -p $VTZ$	207.98	0.94	1.24	1.0967	$-0.0020$	$-0.0024$	2423.8	7.9	10.5
	$cc$ -p $VQZ$	214.17	0.94	1.25	1.0931	$-0.0020$	$-0.0024$	2435.4	7.9	10.6
CCSD(T)	$cc$ -p $VDZ$	200.69	0.79	1.07	1.1189	$-0.0019$	$-0.0023$	2339.1	7.5	10.3
	$cc$ - $pVTZ$	216.70	0.88	1.18	1.1038	$-0.0020$	$-0.0025$	2346.0	8.1	10.8
	$cc-pVQZ$	223.17	0.88	1.18	1.1003	$-0.0020$	$-0.0025$	2356.2	8.1	10.8

 $^a$  Ref. [6]

bound at the SCF level of theory. Very little differences are also observed between CCSD and CCSD(T) for  $\Delta D_e$ ,  $\Delta r_e$ , and  $\Delta \omega_e$ , implying that only modest amounts of valence electron correlation need be included when using CPPs to estimate core-valence correlation contributions. Therefore it should be stressed that in assessing the accuracy of the CPP results, they should be compared to the best estimates for the core-valence correlation effect

Method	Basis set	$D_e$ (kcal/mol)			$r_e$ (Å)			$\omega_e$ (cm <sup>-1</sup> )			
		Valence	CPP/1	CPP/2	Valence	CPP/1	CPP/2	Valence	CPP/1	CPP/2	
$CCSD(T)/cc$ -p $CV6Z^a$		119.16		$+0.22$	1.2063		$-0.0021$	1604.1	$+6.0$		
Best estimate <sup>a</sup>				$+0.3$		$-0.002$			$+6$		
<b>SCF</b>	$cc$ -p $VDZ$ $cc$ -p $VTZ$ $cc-pVQZ$	25.87 32.80 34.00	0.44 0.28 0.27	0.61 0.43 0.42	1.1543 1.1523 1.1510	$-0.0016$ $-0.0017$ $-0.0017$	$-0.0021$ $-0.0022$ $-0.0021$	2021.1 2004.2 2007.8	6.3 5.3 5.1	8.5 7.4 7.1	
<b>CCSD</b>	$cc$ -p $VDZ$ $cc$ -p $VTZ$ $cc$ -p $VQZ$	98.03 105.14 108.29	0.46 0.30 0.31	0.62 0.45 0.46	1.2054 1.1993 1.1947	$-0.0017$ $-0.0017$ $-0.0016$	$-0.0022$ $-0.0022$ $-0.0021$	1670.1 1677.2 1694.8	5.4 3.8 3.8	7.4 5.7 5.6	
CCSD(T)	$cc$ -p $VDZ$ $cc$ -p $VTZ$ $cc$ -p $VQZ$	104.03 113.63 117.19	0.41 0.26 0.26	0.55 0.39 0.39	1.2158 1.2120 1.2077	$-0.0018$ $-0.0018$ $-0.0017$	$-0.0024$ $-0.0023$ $-0.0022$	1587.7 1585.4 1600.1	5.7 3.8 3.8	7.9 5.7 5.7	

**Table 5.** Calculated spectroscopic constants of  $O<sub>2</sub>$  for valence-only correlation and the changes in these values due to using core polarization potentials

 $a$  Ref. [6]

**Table 6.** Calculated spectroscopic constants of  $F<sub>2</sub>$  for valence-only correlation and the changes in these values due to using core polarization potentials

Method	Basis set	$D_e$ (kcal/mol)			$r_e$ (Å)			$\omega_e$ (cm <sup>-1</sup> )		
		Valence	CPP/1	CPP/2	Valence	CPP/1	CPP/2	Valence	CPP/1	CPP/2
$CCSD(T)/cc$ -p $CV6Z^a$ Best estimate <sup>a</sup>		38.26	$-0.1$	$-0.07$	1.4102		$-0.0015$ $-0.0015$	928.4	$+1.6$ $+1$	
<b>SCF</b>	$cc$ -p $VDZ$	$-33.40$	0.04	0.05	1.3477	$-0.0018$	$-0.0023$	1180.2	2.5	3.1
	$cc$ -p $VTZ$	$-25.32$	$-0.06$	$-0.04$	1.3291	$-0.0012$	$-0.0016$	1266.9	1.2	1.8
	$cc-pVQZ$	$-24.78$	$-0.05$	$-0.04$	1.3275	$-0.0012$	$-0.0016$	1263.9	1.4	1.9
	aug-cc-pVDZ	$-30.54$	$-0.01$	0.01	1.3382	$-0.0016$	$-0.0020$	1215.8	2.1	2.7
	aug-cc-pVTZ	$-24.64$	$-0.06$	$-0.04$	1.3282	$-0.0012$	$-0.0016$	1270.8	1.3	1.8
	aug-cc-pVQZ	$-24.48$	$-0.05$	$-0.04$	1.3267	$-0.0012$	$-0.0016$	1265.9	1.4	1.9
<b>CCSD</b>	$cc$ -p $VDZ$	22.28	0.06	0.07	1.4321	$-0.0017$	$-0.0021$	885.9	1.0	1.2
	$cc$ -p $VTZ$	28.38	0.00	0.01	1.3946	$-0.0010$	$-0.0013$	1012.4	0.2	0.4
	$cc-pVQZ$	29.89	0.01	0.02	1.3907	$-0.0010$	$-0.0013$	1015.5	0.2	0.4
	aug-cc-pVDZ	24.58	0.02	0.03	1.4255	$-0.0014$	$-0.0018$	921.2	0.5	0.7
	aug-cc-pVTZ	29.80	0.00	0.02	1.3957	$-0.0010$	$-0.0013$	1010.7	0.3	0.5
	$aug-cc-pVQZ$	30.77	0.01	0.02	1.3903	$-0.0010$	$-0.0013$	1016.6	0.3	0.5
CCSD(T)	$cc$ -p $VDZ$	27.24	0.02	0.02	1.4576	$-0.0019$	$-0.0023$	784.0	1.2	1.5
	$cc$ -p $VTZ$	34.98	$-0.02$	$-0.02$	1.4158	$-0.0010$	$-0.0013$	919.9	0.2	0.5
	$cc-pVQZ$	36.92	$-0.02$	$-0.01$	1.4130	$-0.0010$	$-0.0014$	921.1	0.2	0.4
	aug-cc-pVDZ	30.03	$-0.01$	$-0.01$	1.4500	$-0.0015$	$-0.0020$	826.2	0.6	0.9
	aug-cc-pVTZ	36.65	$-0.02$	$-0.02$	1.4181	$-0.0010$	$-0.0014$	916.3	0.3	0.5
	aug-cc-pVQZ	37.90	$-0.02$	$-0.01$	1.4130	$-0.0010$	$-0.0014$	921.4	0.3	0.5

 $a$  Ref. [6]

and not, for instance, to just large basis set CCSD(T) results when CCSD(T) is used to correlate the valence electrons in the CPP calculations.

The second trend to be observed is the relative insensitivity of the CPP results to the basis set. It is generally observed in Tables 2–6 that the core-valence effects obtained by the CPPs are nearly completely converged with only the cc-pVTZ basis set. In the case of  $F_2$  the effects of additional diffuse functions were also investigated by using the aug-cc-pVnZ basis sets. Improvements were observed at the double zeta level, but the aug-cc-pVTZ results were essentially identical to the values obtained with the standard cc-pVTZ basis set. Both of these trends are in strong contrast to the conventional method of predicting core-valence correlation effects on spectroscopic constants, where large basis sets and sophisticated correlation methods are required to obtain reliable values.

Upon comparing the core-valence correlation effects predicted by the CPP calculations with our "best estimates" from the literature (Tables 2–6 and Fig. 1), it is observed that CPP/2 is closer to the "best estimate" values compared to CPP/1 for  $B_2$ ,  $C_2$ , and  $F_2$ . For these three species, the CPP/2 values differ from the "best" estimates'' by an average of less than 0.1 kcal/mol for  $\Delta D_e$ , less than a few tenths of a mÅ in  $\Delta r_e$ , and less than  $2 \text{ cm}^{-1}$  in  $\Delta \omega_e$ . These are certainly well within the uncertainty of the best estimate values. In particular the large (positive) value of  $\Delta D_e$  for  $C_2$  is very well reproduced by CPP/2, as well as the small negative value for  $F<sub>2</sub>$ . For each of these species, CPP/1 underestimates the core-valence correlation effects.

Method	Basis set	$D_e$ (kcal/mol)			$r_e$ (Å)			$\omega_e$ (cm <sup>-1</sup> )		
		Valence	CPP/1	CPP/2	Valence	CPP/1	CPP/2	Valence	CPP/1	CPP/2
$CCSD(T)/cc$ -p $CV5Z^a$ Best estimate <sup>a</sup>		257.69	$+0.93$ $+0.95$		1.1306	$-0.0024$ $-0.0026$		2165.3		$+9.8$ $+10$
<b>SCF</b>	cc-pVDZ	176.69	0.79	1.04	1.1101	$-0.0022$	$-0.0026$	2431.6	7.5	10.2
	$cc$ -p $VTZ$	183.28	0.89	1.19	1.1045	$-0.0024$	$-0.0029$	2425.0	8.9	11.7
	cc-pVQZ	184.75	0.91	1.21	1.1020	$-0.0024$	$-0.0029$	2427.3	9.1	12.0
	aug-cc-pCVDZ	178.21	0.78	1.04	1.1099	$-0.0023$	$-0.0028$	2401.8	8.5	11.3
	aug-cc-pCVTZ	183.97	0.89	1.18	1.1037	$-0.0024$	$-0.0029$	2421.0	9.0	11.9
	aug-cc-pCVQZ	184.91	0.91	1.20	1.1019	$-0.0024$	$-0.0029$	2426.5	9.1	12.0
<b>CCSD</b>	$cc-pVDZ$	235.73	0.91	1.21	1.1384	$-0.0021$	$-0.0026$	2209.2	7.0	9.6
	cc-pVTZ	244.46	0.97	1.30	1.1288	$-0.0023$	$-0.0028$	2222.9	8.1	10.9
	$cc-pVQZ$	248.67	0.99	1.31	1.1243	$-0.0023$	$-0.0028$	2234.8	8.2	11.0
	aug-cc-pCVDZ	235.52	0.92	1.23	1.1390	$-0.0023$	$-0.0028$	2171.7	8.0	10.8
	aug-cc-pCVTZ	245.04	0.96	1.29	1.1281	$-0.0023$	$-0.0029$	2216.7	8.1	11.0
	aug-cc-pCVQZ	248.96	0.98	1.30	1.1243	$-0.0023$	$-0.0028$	2231.7	8.2	11.1
CCSD(T)	$cc$ -p $VDZ$	241.52	0.87	1.17	1.1446	$-0.0022$	$-0.0027$	2143.9	7.2	9.9
	$cc$ -p $VTZ$	251.90	0.93	1.25	1.1357	$-0.0024$	$-0.0029$	2153.7	8.2	11.0
	$cc-pVQZ$	256.32	0.94	1.26	1.1314	$-0.0023$	$-0.0029$	2164.4	8.3	11.2
	aug-cc-pCVDZ	241.61	0.88	1.18	1.1457	$-0.0023$	$-0.0029$	2104.4	8.1	11.0
	aug-cc-pCVTZ	252.53	0.92	1.24	1.1353	$-0.0024$	$-0.0029$	2145.7	8.2	11.1
	$aug-cc-pCVQZ$	256.65	0.93	1.25	1.1315	$-0.0024$	$-0.0029$	2160.4	8.3	11.2

Table 7. Calculated spectroscopic constants of CO for valence-only correlation and the changes in these values due to using core polarization potentials

<sup>a</sup> Ref. [42]. The ACPF/cc-pCV5Z core-valence correlation effects from [42] were +0.90 kcal/mol, -0.0024 Å, and +9.4 cm<sup>-1</sup>

Table 8. Calculated spectroscopic constants of CN for valence-only correlation and the changes in these values due to using core polarization potentials

Method	Basis set	$D_e$ (kcal/mol)			$r_e(A)$			$\omega_e$ (cm <sup>-1</sup> )		
		Valence	CPP/1	CPP/2	Valence	CPP/1	CPP/2	Valence	CPP/1	CPP/2
	$ACPF$ /cc-p $CV5Z^a$		1.15		1.1754	$-0.0028$		2056.9	$+9.5$	
<b>SCF</b>	$cc$ -p $VDZ$	80.28	0.77	1.18	1.1390	$-0.0023$	$-0.0028$	2460.9	9.1	12.8
	$cc$ -p $VTZ$	86.47	0.93	1.39	1.1292	$-0.0025$	$-0.0031$	2450.7	10.2	14.0
	$cc-pVQZ$	87.88	0.93	1.40	1.1273	$-0.0025$	$-0.0032$	2453.3	10.1	13.9
<b>CCSD</b>	$cc$ -p $VDZ$	150.05	0.95	1.36	1.1850	$-0.0023$	$-0.0029$	2115.8	7.8	10.9
	$cc$ -p $VTZ$	160.66	1.06	1.50	1.1683	$-0.0025$	$-0.0031$	2138.3	8.7	11.9
	$cc$ -p $VQZ$	165.57	1.07	1.52	1.1636	$-0.0025$	$-0.0031$	2153.5	8.5	11.8
CCSD(T)	$cc$ -p $VDZ$	158.17	0.86	1.25	1.1948	$-0.0024$	$-0.0030$	2031.2	8.0	11.1
	$cc$ -p $VTZ$	170.43	0.98	1.39	1.1789	$-0.0026$	$-0.0032$	2050.3	8.7	11.8
	$cc$ -p $VQZ$	175.61	0.99	1.41	1.1744	$-0.0025$	$-0.0032$	2064.3	8.6	11.8

<sup>a</sup>This work (best estimate). Pradhan et al. [47] also reported a value for  $\Delta D_e$  of 1.18 kcal/mol using ACPF wave functions

Table 9. Calculated spectroscopic constants of CH for valence-only correlation and the changes in these values due to using core polarization potentials

Method	Basis set	$D_e$ (kcal/mol)			$r_e(A)$			$\omega_e$ (cm <sup>-1</sup> )		
		Valence	CPP/1	CPP/2	Valence	CPP/1	CPP/2	Valence	CPP/1	CPP/2
$CCSD(T)/cc$ -p $CVQZ^a$ $ACPF$ $\c{c}$ - $pC\overline{V}5Z^b$		83.08 83.41	$+0.13$ $+0.22$		1.1202 1.1201		$-0.0018$ $-0.0021$	2854.1 2851.5		$+8.0$ $+10.5$
<b>SCF</b>	$cc$ -p $VDZ$	54.65	0.02	0.08	1.1179	$-0.0015$	$-0.0019$	3048.7	5.6	8.7
	$cc$ -p $VTZ$	56.76	0.05	0.12	1.1053	$-0.0017$	$-0.0022$	3036.4	6.6	9.8
	$cc-pVQZ$	57.13	0.04	0.12	1.1039	$-0.0017$	$-0.0022$	3043.5	6.9	10.1
<b>CCSD</b>	$cc$ -p $VDZ$	74.95	0.08	0.15	1.1418	$-0.0015$	$-0.0020$	2831.6	5.1	8.3
	$cc$ -p $VTZ$	80.72	0.12	0.20	1.1207	$-0.0017$	$-0.0022$	2859.0	6.4	9.6
	$cc$ -p $VQZ$	82.24	0.11	0.19	1.1185	$-0.0017$	$-0.0022$	2873.8	6.4	9.6
CCSD(T)	$cc$ -p $VDZ$	75.53	0.08	0.14	1.1431	$-0.0015$	$-0.0020$	2816.7	5.3	8.6
	$cc$ -p $VTZ$	81.50	0.11	0.20	1.1224	$-0.0017$	$-0.0023$	2840.6	6.5	9.7
	$cc-pVQZ$	83.07	0.11	0.19	1.1202	$-0.0017$	$-0.0023$	2854.3	6.5	9.8

<sup>a</sup>Ref. [43]. Partridge and Bauschlicher [46] have reported a values of +0.17 kcal/mol and -0.0017 Å for  $\Delta D_e$  and  $\Delta r_e$ , respectively, using  $CCSD(T)$  wave functions<br>  ${}^{b}This$  work (best estimate)

Table 10. Calculated spectroscopic constants of HF for valence-only correlation and the changes in these values due to using core polarization potentials

Method	Basis set	$D_e$ (kcal/mol)			$r_e$ (Å)			$\omega_e$ (cm <sup>-1</sup> )		
		Valence	CPP/1	CPP/2	Valence	CPP/1	CPP/2	Valence	CPP/1	CPP/2
$CCSD(T)/cc$ -p $CV5Z^a$		141.03		$+0.18$	0.9168	$-0.0006$		4150.0	$+4.6$	
$ACPF$ /cc-p $CV5Z^b$		139.89		$+0.23$	0.9163		$-0.0007$	4153.7		$+6.0$
<b>SCF</b>	$cc$ -p $VDZ$	93.72	0.13	0.20	0.9015	$-0.0006$	$-0.0007$	4440.2	4.6	6.1
	$cc$ - $pVTZ$	100.06	0.13	0.20	0.8979	$-0.0006$	$-0.0007$	4481.1	4.4	5.9
	$cc$ -p $VQZ$	101.19	0.13	0.20	0.8969	$-0.0006$	$-0.0007$	4476.5	4.2	5.7
	aug-cc-pVDZ	99.48	0.13	0.20	0.9002	$-0.0005$	$-0.0007$	4466.3	4.6	6.1
	aug-cc-pVTZ	101.37	0.12	0.20	0.8991	$-0.0006$	$-0.0007$	4464.8	4.3	5.8
	aug-cc-pVQZ	101.61	0.13	0.20	0.8973	$-0.0006$	$-0.0007$	4471.0	4.2	5.7
<b>CCSD</b>	$cc-pVDZ$	125.73	0.15	0.22	0.9190	$-0.0005$	$-0.0007$	4168.6	4.1	5.6
	cc-pVTZ	135.32	0.15	0.22	0.9152	$-0.0005$	$-0.0007$	4211.5	3.6	5.1
	$cc-pVQZ$	138.10	0.15	0.23	0.9137	$-0.0005$	$-0.0007$	4204.2	3.5	5.0
	aug-cc-pVDZ	133.40	0.15	0.22	0.9222	$-0.0006$	$-0.0007$	4115.9	4.2	5.7
	aug-cc-pVTZ	137.45	0.14	0.22	0.9182	$-0.0005$	$-0.0007$	4169.3	3.5	4.9
	aug-cc-pVQZ	138.98	0.15	0.22	0.9149	$-0.0005$	$-0.0007$	4188.6	3.5	5.0
CCSD(T)	$cc$ -p $VDZ$	126.37	0.14	0.22	0.9199	$-0.0005$	$-0.0007$	4149.8	4.2	5.8
	$cc$ -p $VTZ$	137.03	0.14	0.22	0.9173	$-0.0006$	$-0.0007$	4176.9	3.7	5.2
	$cc-pVQZ$	140.06	0.14	0.22	0.9162	$-0.0005$	$-0.0007$	4161.8	3.6	5.1
	aug-cc-pVDZ	134.52	0.14	0.21	0.9241	$-0.0006$	$-0.0007$	4080.3	4.3	5.9
	aug-cc-pVTZ	139.40	0.14	0.21	0.9210	$-0.0006$	$-0.0007$	4124.3	3.6	5.0
	aug-cc-pVQZ	141.04	0.14	0.22	0.9177	$-0.0005$	$-0.0007$	4141.3	3.6	5.1

 $\sum_{b}^{a}$  Ref. [41]<br> $\sum_{b}$ This work (best estimate)

Method	Basis set	$D_e$ (HC-CH) (kcal/mol)			$r_{\rm e}$ (CC) (Å)			$r_{\rm e}$ (CH) (Å)		
		Valence	CPP/1	CPP/2	Valence	CPP/1	CPP/2	Valence	CPP/1	CPP/2
CCSD(T)/ACPF <sup>a</sup>			2.14 / 2.35							
<b>SCF</b>	$cc$ -p $VDZ$	181.05	1.40	2.04	1.1918	$-0.0024$	$-0.0029$	1.0638	$-0.0012$	$-0.0014$
	$cc$ -p $VTZ$	186.22	1.61	2.32	1.1801	$-0.0027$	$-0.0033$	1.0540	$-0.0014$	$-0.0016$
	$cc-pVQZ$	186.56	1.65	2.36	1.1792	$-0.0028$	$-0.0033$	1.0540	$-0.0014$	$-0.0016$
<b>CCSD</b>	$cc$ -p $VDZ$	215.38	1.56	2.27	1.2227	$-0.0023$	$-0.0028$	1.0775	$-0.0012$	$-0.0014$
	$cc$ -p $VTZ$	224.39	1.70	2.45	1.2032	$-0.0025$	$-0.0031$	1.0620	$-0.0013$	$-0.0015$
	$cc-pVQZ$	227.21	1.73	2.47	1.1999	$-0.0025$	$-0.0030$	1.0616	$-0.0013$	$-0.0015$
CCSD(T)	$cc-pVDZ$	220.27	1.52	2.23	1.2287	$-0.0024$	$-0.0028$	1.0789	$-0.0012$	$-0.0014$
	$cc$ -p $VTZ$	230.46	1.66	2.41	1.2096	$-0.0026$	$-0.0031$	1.0637	$-0.0013$	$-0.0015$
	$cc-pVQZ$	233.51	1.68	2.43	1.2064	$-0.0026$	$-0.0031$	1.0633	$-0.0013$	$-0.0015$

**Table 11.** Calculated spectroscopic constants of  $C_2H_2$  for valence-only correlation and the changes in these values due to using core polarization potentials

 $a$  Ref. [46].

In the case of  $N_2$ , and to a lesser extent  $O_2$ , the corevalence correlation effect is overestimated by CPP/2. For  $N_2$  the CCSD(T)/cc-pVTZ results using CPP/2 are 1.18 kcal/mol,  $-0.0025$  Å, and  $+10.8$  cm<sup>-1</sup> for  $\Delta D_e$ ,  $\Delta r_e$ , and  $\Delta\omega_e$ , respectively. These can be compared to best estimates of  $+0.9$  kcal/mol,  $-0.002$  Å, and  $+9$  cm<sup>-1</sup>. For  $O_2$  the largest discrepancy is the value of  $\Delta D_e$  predicted by CPP/2, which is probably too large by about 0.1 kcal/mol. The core-valence contributions to  $r_e$  and  $\omega_e$ , however, are in good agreement for CPP/2. The source of these errors, which also occur in CO and CN (see below), is not yet understood, but the restriction to frozen, spherical cores in the CPP calculations could perhaps give rise to these overestimations. Work is currently in progress to remove this restriction. Other forms of the cutoff function are also being investigated. Given the large computational expense, however, to accurately calculate core-valence correlation effects using standard correlation methods, errors of a few tenths of a kcal/mol would appear to be very acceptable when there are essentially no additional costs beyond the usual valence-only calculation.

## 3.2 CO and CN

The results for CO and CN given in Tables 7 and 8 show very similar trends as were discussed above for the homonuclear diatomics. In the case of CO, very similar observations as noted for  $N_2$  can be made. For example,

the value of  $\Delta D_e$  obtained using CPP/2 is too large by about 0.3 kcal/mol out of 1 kcal/mol. The effects of using basis sets that include tight  $d$  and  $f$  functions designed to recover core-core and core-valence correlation effects were also investigated for CO. The results of these calculations, which actually used the aug-ccpCVDZ through aug-cc-pCVQZ basis sets (extra diffuse and tight functions), are also shown in Table 7 and, as with the diffuse augmented sets alone, result in very few benefits. The results for CN are intermediate in accuracy between  $C_2$  and CO, e.g.,  $\Delta D_e$  is too large by about 0.2 kcal/mol and  $\Delta r_e$  is too negative by about 0.0005 Å).

## 3.3 CH and HF

For both CH and HF, CPP/2 yields very accurate values for the core-valence contributions to  $D_e$ ,  $r_e$ , and  $\omega_e$ (Tables 9 and 10). In the all-electron benchmark calculations, the previous CCSD(T) results [43, 46] for CH yielded core-valence correlation effects somewhat smaller than those obtained by the ACPF method. The values obtained from the latter calculations are in very good agreement with the ones from CPP/2, i.e., within 0.03 kcal/mol, 0.0002 Å, and 1 cm<sup>-1</sup> for  $\Delta D_e$ ,  $\Delta r_e$ , and  $\Delta\omega_e$ , respectively. Similar excellent agreement between CPP/2 and our best estimates are also observed for HF. In this case, it has been observed previously [59, 60] that extra diffuse functions were very important for obtaining accurate dissociation energies and harmonic frequencies. However, as also shown above for  $F_2$ , the addition of these functions to the basis set have a negligible influence on the core-valence correlation effect predicted by the CPPs.

#### 3.4  $C_2H_2$

The last example considered in the present work was the acetylene molecule. Initially this was chosen in an attempt to show that the short, triple bond in  $N_2$  (and CO) was responsible for the overestimation of the corevalence effects by  $\text{CPP}/2$  in that molecule. The results shown in Table 11 for  $C_2H_2$ , however, indicate that this is probably not a valid explanation since excellent results are obtained by  $\text{CPP}/2$ . The core-valence effect on  $D_e(HC-CH)$  is surprisingly large, as first shown by Partridge and Bauschlicher [46]. Use of CPP/2 reproduces their ACPF result  $(+2.35 \text{ kcal/mol})$  to within 0.1 kcal/mol. In contrast, CPP/1 strongly underestimates the core-valence contributions. Our calculations also yield core-valence contributions to the two bond lengths in  $C_2H_2$ , predicting bond length contractions (CPP/2) of 0.0031 and 0.0015 A for  $r_e$ (CC) and  $r_e$ (CH), respectively. These should be accurate to within  $0.2$  mA.

## 4 Conclusions

The core-valence correlation contributions to the spectroscopic constants  $(D_e, r_e, \omega_e)$  of several molecules containing first-row atoms have been investigated using effective CPPs. In each case the CPP results reproduced well the best estimates of these quantities, which were obtained in extensive calculations where all of the electrons were included in the correlation treatment. The use of the CPP method reduced the error in the calculated valence  $D_e$  values (relative to the all-electron  $D_e$  values) from a maximum of nearly 2.5 kcal/mol to just 0.3 kcal/mol. Similar error reductions were found for  $r_e$  and  $\omega_e$ .

As shown in Fig. 1, this study indicates that the choice of the cut-off function in the CPP plays a more important role than anticipated by Meyer and coworkers [12]. In contrast to their claim that their different cut-off functions performed equally well as long as the parameters were carefully adjusted, we found CPP/2 to outperform CPP/1 in nearly all cases. This was especially true when the core-valence effects became sizeable. However, it should also be mentioned that the fitting procedure devised by Meyer is obviously better suited for alkali and alkaline earth elements. For elements to the right of the periodic table, the one valenceelectron ions used in the fitting procedure for the cut-off function have little similarity to the atoms in the neutral molecules of the present study. So it is somewhat surprising that one of the cut-off functions  $(CPP/2)$  does so well, while  $CPP/1$  exhibits significant deviations. Gaining more insight into the role of the cut-off function is the topic of an ongoing investigation.

In general, the minimal computational expense and apparent high accuracy of the CPPs in reproducing the core-valence correlation effects should result in a more widespread and routine use of these operators in molecular calculations.

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